

# ON INFINITE SERIES OF THREE VARIABLES INVOLVING WHITTAKER AND BESSEL FUNCTIONS

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**Abstract.** The main object of the present paper is to obtain on infinite series of three variables involving Whittaker and Bessel function into a series of Srivastava's triple hypergeometric function  $F^{(3)}$ . A number of known and knew transformations are also discussed as the special cases of our main result.

## 1 Introduction

Bessel function of order  $\nu$  defined by Srivasta and Manocha as follows [10]:

$$J_\nu(z) = \frac{(z/2)^\nu}{\Gamma(\nu + 1)} {}_0F_1 \left[ \begin{matrix} -; \\ \nu + 1; \end{matrix} \frac{-z^2}{4} \right]. \tag{1.1}$$

Modified Bessel function defined by Rainville [7] as follows:

$$I_\nu(z) = \frac{(z/2)^\nu}{\Gamma(\nu + 1)} {}_0F_1 \left[ \begin{matrix} -; \\ \nu + 1; \end{matrix} \frac{z^2}{4} \right]. \tag{1.2}$$

Bessel functions are of most frequent use in the theory of integral transform. The modified Bessel function  $I_\nu(z)$  is related to  $J_\nu(z)$  in much the same way that the hyperbolic function is related to trigonometric function, and we have

$$I_\nu(z) = i^{-\nu} J_\nu(iz). \tag{1.3}$$

Transformation of Appell's function  $F_4$  into a series of Srivastava's triple hypergeometric function  $F^{(3)}$  given by Khan and Pathan [6] as follows:

$$\begin{aligned} & F_4 \left[ \begin{matrix} a, a \\ \frac{a}{2}, \frac{a}{2} \end{matrix}; \frac{1}{2}; b, c; \frac{4x}{(1+x+y)^2}, \frac{y^2}{(1+x+y)^2} \right] \\ &= \frac{(1+x+y)^{2a} \Gamma(a + \frac{1}{2}) \Gamma(a - b + \frac{3}{2})}{\sqrt{\pi} \Gamma(2a - b + \frac{3}{2})} \\ & \cdot \left\{ F^{(3)} \left[ \begin{matrix} a, a + \frac{1}{2} :: a - \frac{1}{2}b + \frac{3}{4}; a - \frac{1}{2}b + \frac{5}{4}; a - b + 1; a - b + \frac{3}{2}; -; -; \\ - : - : c - \frac{1}{2}; -; \frac{b}{2} - \frac{1}{4}, \frac{b}{2} + \frac{1}{4} 2c - 1, a - b + 1; b, \frac{1}{2} \end{matrix} \right. \right. \\ & \qquad \qquad \qquad \left. \left. -2y, x, 1 \right] \right. \\ & \qquad \qquad \qquad \left. \frac{a(b - \frac{1}{2})}{(a - \frac{1}{2}b + \frac{3}{4})} \right. \\ & \cdot F^{(3)} \left[ \begin{matrix} a + \frac{1}{2}, a + 1, a - \frac{1}{2}b + \frac{5}{4}, a - \frac{1}{2}b + \frac{7}{4}; a - b + 1; a - b + \frac{3}{2}; -; -; \\ - : - : -; c - \frac{1}{2}; -; \frac{b}{2} + \frac{1}{4}, \frac{b}{2} + \frac{3}{4}; 2c - 1, a - b + 1; b, \frac{3}{2} \end{matrix} \right. \\ & \qquad \qquad \qquad \left. \left. -2y, x, 1 \right] \right\}, \tag{1.4} \end{aligned}$$

where  $F^{(3)}[x, y, z]$  is Srivastava's triple hypergeometric series defined by Srivastava and Manocha [10, p.69(39)] as follows:

$$\begin{aligned}
 & F^{(3)} \left[ \begin{matrix} (a) :: (b); (b'); (b''); (e); (e'); (e'') & ; \\ (d) :: (g); (g'); (g''); (h); (h'); (h'') & ; \end{matrix} \right. \\
 & \qquad \qquad \qquad \left. x, y, z \right] \\
 &= \sum_{m,n,p=0}^{\infty} \frac{[(a)]_{m+n+p} [(b)]_{m+n} [(b')]_{n+p} [(b'')]_{p+m} [(e)]_m [(e')]_n [(e'')]_p}{[(d)]_{m+n+p} [(g)]_{m+n} [(g'')]_{n+p} [(g'')]_{p+m} [(h)]_m [(h')]_n [(h'')]_p} \frac{x^m y^n z^p}{m! n! p!}. \tag{1.5}
 \end{aligned}$$

**2 Main Transformation**

In order to obtain the main transformation of the present paper, we establish an integral in the following form:

$$\begin{aligned}
 & \int_0^{\infty} t^{\sigma-1} e^{-zt} W_{k,m}(ht) J_{\nu}(\alpha t) {}_1F_2 \left( \begin{matrix} \beta & ; \\ \gamma, \delta & ; \end{matrix} \middle| y^2 t^2 \right) dt \\
 &= \frac{h^{m+\frac{1}{2}} \alpha^{\nu} \Gamma(C) \Gamma(C')}{2^{\nu} \Gamma(\nu+1) \Gamma(E) (z+h/2)^C} \\
 & \cdot \left\{ F^{(3)} \left[ \begin{matrix} \frac{C}{2}, \frac{C+1}{2} :: \frac{C'}{2}, \frac{C'+1}{2}; -; -; \beta; -; \frac{D}{2}, \frac{D+1}{2}; & \frac{4y^2}{(z+h/2)^2}, \frac{-\alpha^2}{(z+h/2)^2}, \left(\frac{z-h/2}{z+h/2}\right)^2 \\ \frac{E}{2}, \frac{E+1}{2} :: -; -; -; \gamma, \delta; \nu+1; \frac{1}{2} & ; \end{matrix} \right] \right. \\
 & \qquad \qquad \qquad \left. + \frac{CD}{E} \left(\frac{z-h/2}{z+h/2}\right) \right. \\
 & \cdot F^{(3)} \left[ \begin{matrix} \frac{C+1}{2}, \frac{C+2}{2} :: \frac{C'}{2}, \frac{C'+1}{2}; -; -; \beta; -; \frac{D+1}{2}, \frac{D+2}{2}; & \frac{4y^2}{(z+h/2)^2}, \frac{-\alpha^2}{(z+h/2)^2}, \left(\frac{z-h/2}{z+h/2}\right)^2 \\ \frac{E+1}{2}, \frac{E+2}{2} :: -; -; -; \gamma, \delta; \nu+1; \frac{3}{2} & ; \end{matrix} \right] \left. \right\}, \tag{2.1}
 \end{aligned}$$

where  $C = \sigma + \nu + m + \frac{1}{2}$ ,  $C' = \sigma + \nu - m + \frac{1}{2}$ ,  $D = m - k + \frac{1}{2}$ ,  $E = \sigma + \nu - k + 1$ ,  $\text{Re}(C) > 0$ ,  $\text{Re}(C') > 0$ ,  $\text{Re}(z+h/2) > 0$  and

$$\begin{aligned}
 W_{k,m}(x) &= \frac{\Gamma(-2m)}{\Gamma(\frac{1}{2} - m - k)} M_{k,m}(x) + \frac{\Gamma(2m)}{\Gamma(\frac{1}{2} + m - k)} M_{k,-m}(x). \\
 M_{k,m}(x) &= x^{m+1/2} e^{(-1/2)x} {}_1F_1 \left( \frac{1}{2} + m - k, 2m + 1; x \right).
 \end{aligned}$$

To obtain equation (2.1), express  $J_{\nu}(\alpha t)$  in terms of  ${}_0F_1$  with the help of the equation (1.1) and then expanding  ${}_0F_1$  and  ${}_1F_2$  in series and using the integral transform [2, p.216(16)], we get

$$\begin{aligned}
 I &= \frac{(\alpha/2)^{\nu} h^{m+1/2}}{\Gamma(\nu+1)} \sum_{p=0}^{\infty} \sum_{q=0}^{\infty} \frac{(-\alpha^2)^p (\beta)_q y^{2q}}{(\nu+1)_p (\gamma)_q (\delta)_q p! q!} \\
 & \cdot \frac{\Gamma(\sigma + \nu \pm m + \frac{1}{2} + 2p + 2q)}{\Gamma(\sigma + \nu - k + 1 + 2p + 2q) (z+h/2)^{\sigma+\nu+m+\frac{1}{2}+2p+2q}} \\
 & \cdot {}_2F_1 \left( \begin{matrix} \sigma + \nu + m + \frac{1}{2} + 2p + 2q, m - k + \frac{1}{2}; & \frac{z-h/2}{z+h/2} \\ \sigma + \nu - k + 1 + 2p + 2q & ; \end{matrix} \right).
 \end{aligned}$$

Now making use of a following result of Carlson [4, p.234(10)], in the above equation

$${}_2F_1 \left( \begin{matrix} a, b; & \\ c & ; \end{matrix} \middle| x \right) = {}_4F_3 \left( \begin{matrix} \frac{a}{2}, \frac{a+1}{2}, \frac{b}{2}, \frac{b+1}{2}; & \\ \frac{1}{2}, \frac{c}{2}, \frac{c+1}{2} & ; \end{matrix} \middle| x^2 \right) + \frac{ab}{c} x {}_4F_3 \left( \begin{matrix} \frac{a+1}{2}, \frac{a+2}{2}, \frac{b+1}{2}, \frac{b+2}{2}; & \\ \frac{3}{2}, \frac{c+1}{2}, \frac{c+2}{2} & ; \end{matrix} \middle| x^2 \right), \tag{2.2}$$

and then expanding  ${}_4F_3$ 's in series and interpreting in the form of  $F^{(3)}$ , we arrive at the main integral (2.1).

Taking  $k = 0, h = 2x, m = \nu, \beta = \mu + \frac{1}{2}, \gamma = \lambda = 2\mu + 1, \delta = \lambda + \mu$ , replacing  $\alpha$  by  $i\alpha$ ,  $y$  by  $iy$  and using equation (1.3) and a result [3, p.432] in equation (2.1) and comparing it with a result of Saxena [11, p.131(9)], we get

$$\int_0^\infty t^{2\lambda+2\mu-1} e^{-zt} K_\nu(xt) I_\nu(\alpha t) {}_1F_2 \left( \begin{matrix} \mu + \frac{1}{2} & ; \\ & -y^2 t^2 \end{matrix} \right) dt$$

$$= \frac{2^{2(\lambda+\mu-1)} (\alpha\beta)^\nu \Gamma(\lambda + \mu + \nu) \Gamma(\lambda + \mu)}{(\alpha^2 + \beta^2 + 2\gamma^2)^{\lambda+\mu+\nu} \Gamma(\nu + 1)}$$

$$\cdot F_4 \left[ \frac{\lambda + \mu + \nu}{2}, \frac{\lambda + \mu + \nu + 1}{2}, \nu + 1; \mu + 1; \frac{4\alpha^2 \beta^2}{(\alpha^2 + \beta^2 + 2\gamma^2)^2}, \frac{4\gamma^2}{(\alpha^2 + \beta^2 + 2\gamma^2)^2} \right], \quad (2.3)$$

$\text{Re}(\lambda + \mu + \nu) > 0, \text{Re}(\lambda + \mu) > 0$  and  $\text{Re}(\alpha) > | \text{Re}(\beta) | + | \text{Im}\gamma |$ , then we get the transformation (1.4).

### 3 Special Cases

The following special cases of the main transformation (2.1) are given below:

(i) On setting  $z = h/2$  in equation (2.1), the integral reduces to

$$\int_0^\infty t^{\sigma-1} e^{-(h/2)t} W_{k,m}(ht) J_\nu(\alpha t) {}_1F_2 \left( \begin{matrix} \beta & ; & y^2 t^2 \\ \gamma, \delta & ; & \end{matrix} \right) dt$$

$$= \frac{h^{m+\frac{1}{2}} \alpha^\nu \Gamma(C) \Gamma(C')}{2^\nu \Gamma(\nu + 1) \Gamma(E) h^C} F^{(2)} \left[ \begin{matrix} \frac{C}{2}, \frac{C+1}{2} :: \frac{C'}{2}; \frac{C'+1}{2}; -; -; \beta; -; & \frac{4y^2}{h^2}, \frac{-\alpha^2}{h^2} \\ \frac{E}{2}, \frac{E+1}{2} :: -; \gamma, \delta, \nu + 1; -; & \end{matrix} \right], \quad (3.1)$$

where  $F^{(2)}$  is Kampe' de Ferie't's function defined in Burchnall and Chaundy [1] as follows:

$$F^{(2)} \left[ \begin{matrix} a : b; c & ; \\ x, y & \\ e; f; g & ; \end{matrix} \right] = \sum_{m=0}^\infty \sum_{n=0}^\infty \frac{(a)_{m+n} (b)_m (c)_n x^m y^n}{(e)_{m+n} (f)_m (g)_n m! n!}.$$

(ii) On setting  $k = m + 1/2$  in equation (2.1) and using the relation [4, p.32], then the integral reduces to

$$\int_0^\infty t^{\sigma+m+1/2-1} e^{-(z+h/2)t} J_\nu(\alpha t) {}_1F_2 \left( \begin{matrix} \beta & ; & y^2 t^2 \\ \gamma, \delta & ; & \end{matrix} \right) dt$$

$$= \frac{\alpha^\nu \Gamma(C)}{2^\nu \Gamma(\nu + 1) (z + h/2)^C} F^{(2)} \left[ \begin{matrix} \frac{C}{2}, \frac{C+1}{2} ::: - : \beta & ; -; & \frac{4y^2}{(z + h/2)^2}, \frac{-\alpha^2}{(z + h/2)^2} \\ - :: - : \gamma, \delta, \nu + 1; -; & \end{matrix} \right]. \quad (3.2)$$

(iii) On setting  $k = 0, h = 2x, m = \nu + 1/2, \nu = \mu$  and  $y = 0$  in equation (2.1) and using a relation [4, p.32] and integrating with the help of the integral transform [3, p.198(27)], we get

$$\int_0^\infty t^{\sigma+\mu+2r-1/2} e^{-zt} K_{\nu+1/2}(xt) J_\mu(\alpha t) dt$$

$$= \frac{\alpha^\mu}{2x 2^\mu \Gamma(\mu + 1)} \Gamma(\sigma + \mu - \nu + 2r) \Gamma(\sigma + \mu + \nu + 1 + 2r) \sum_{r=0}^\infty \frac{(-1)^r \alpha^{2r}}{2^{2r} (\mu + 1)_r r!}$$

$$s^{-(\sigma+\mu+2r)} P^{-(\sigma+\mu+2r)}(z/x), \tag{3.3}$$

where  $P_\nu^\mu(z)$  is Legendre function [2, p.122(3)] and  $\text{Re}(\sigma + \mu + \nu + 2r) > -1$ ,  $\text{Re}(\sigma + \mu - \nu + 2r) > 0$  and  $\text{Re}(z) > -\text{Re}(x)$ .

(iv) On setting  $\gamma = \beta + \frac{1}{2}$ ,  $\delta = 2\beta - 1$  in equation (2.1) and using a known result [7, p.595(6)], we get

$$\begin{aligned} & \Gamma(\beta - \frac{1}{2})\Gamma(\beta + \frac{1}{2}) \left(\frac{y}{2}\right)^{2-2\beta} \int_0^\infty t^{\sigma+2-2\beta-1} e^{-zt} W_{k,m}(ht) J_\nu(\alpha t) [I_{\beta-3/2}(yt) I_{\beta-1/2}(yt)] dt \\ &= \frac{h^{m+\frac{1}{2}} \alpha^\nu \Gamma(C) \Gamma(C')}{2^\nu \Gamma(\nu + 1) \Gamma(E) (z + h/2)^C} \\ & \left\{ F^{(3)} \left[ \begin{matrix} \frac{C}{2}, \frac{C+1}{2} :: \frac{C'}{2}, \frac{C'+1}{2}; -; -; \beta; -; \frac{D}{2}, \frac{D+1}{2} & ; & \frac{4y^2}{(z + h/2)^2}, \frac{-\alpha^2}{(z + h/2)^2}, \left(\frac{z - h/2}{z + h/2}\right)^2 \\ \frac{E}{2}, \frac{E+1}{2} :: -; -; -; \beta + \frac{1}{2}, 2\beta - 1; \nu + 1; \frac{1}{2}; & & \end{matrix} \right] \right. \\ & \quad \left. + \frac{CD}{E} \left(\frac{z - h/2}{z + h/2}\right) \right. \\ & \left. F^{(3)} \left[ \begin{matrix} \frac{C+1}{2}, \frac{C+2}{2} :: \frac{C'}{2}, \frac{C'+1}{2}; -; -; \beta; -; \frac{D+1}{2}, \frac{D+2}{2} & ; & \frac{4y^2}{(z + h/2)^2}, \frac{-\alpha^2}{(z + h/2)^2}, \left(\frac{z - h/2}{z + h/2}\right)^2 \\ \frac{E+1}{2}, \frac{E+2}{2} :: -; -; -; \beta + \frac{1}{2}, 2\beta - 1; \nu + 1; \frac{3}{2}; & & \end{matrix} \right] \right\}. \tag{3.4} \end{aligned}$$

(v) On setting  $\beta = \frac{3}{2}$ ,  $\delta = 3 - \gamma$  in equation (2.1) and using a known result [7, p.595(10)], we get

$$\begin{aligned} & \frac{\pi(\gamma - 1)(\gamma - 2)}{\sin \gamma \pi} \int_0^\infty t^{\sigma-1} e^{-zt} W_{k,m}(ht) J_\nu(\alpha t) [I_{1-\gamma}(yt) I_{\gamma-1}(yt) + I_{2-\gamma}(yt) I_{\gamma-2}(yt)] dt \\ &= \frac{h^{m+\frac{1}{2}} \alpha^\nu \Gamma(C) \Gamma(C')}{2^\nu \Gamma(\nu + 1) \Gamma(E) (z + h/2)^C} \\ & \left\{ F^{(3)} \left[ \begin{matrix} \frac{C}{2}, \frac{C+1}{2} :: \frac{C'}{2}, \frac{C'+1}{2}; -; -; \frac{3}{2}; -; \frac{D}{2}, \frac{D+1}{2}; & & \frac{4y^2}{(z + h/2)^2}, \frac{-\alpha^2}{(z + h/2)^2}, \left(\frac{z - h/2}{z + h/2}\right)^2 \\ \frac{E}{2}, \frac{E+1}{2} :: -; -; -; \gamma, 3\gamma - 1; \nu + 1; \frac{1}{2}; & & \end{matrix} \right] \right. \\ & \quad \left. + \frac{CD}{E} \left(\frac{z - h/2}{z + h/2}\right) \right. \\ & \left. F^{(3)} \left[ \begin{matrix} \frac{C+1}{2}, \frac{C+2}{2} :: \frac{C'}{2}, \frac{C'+1}{2}; -; -; \frac{3}{2}; -; \frac{D+1}{2}, \frac{D+2}{2}; & & \frac{4y^2}{(z + h/2)^2}, \frac{-\alpha^2}{(z + h/2)^2}, \left(\frac{z - h/2}{z + h/2}\right)^2 \\ \frac{E+1}{2}, \frac{E+2}{2} :: -; -; -; \gamma, 3\gamma - 1; \nu + 1; \frac{3}{2}; & & \end{matrix} \right] \right\}. \tag{3.5} \end{aligned}$$

(vi) On setting  $\beta = \frac{1}{4}$ ,  $\gamma = \frac{1}{2}$ ,  $\delta = \frac{5}{4}$  in equation (2.1) and using a known result [7, p.598(37)], we get

$$\begin{aligned} & \frac{1}{4} \sqrt{\frac{\pi}{2y}} \int_0^\infty t^{\sigma-3/2} e^{-zt} W_{k,m}(ht) J_\nu(\alpha t) [\text{erf}(\sqrt{2yt}) + \text{erfi}(\sqrt{2yt})] dt \\ &= \frac{h^{m+\frac{1}{2}} \alpha^\nu \Gamma(C) \Gamma(C')}{2^\nu \Gamma(\nu + 1) \Gamma(E) (z + h/2)^C} \\ & \left\{ F^{(3)} \left[ \begin{matrix} \frac{C}{2}, \frac{C+1}{2} :: \frac{C'}{2}, \frac{C'+1}{2}; -; -; \frac{1}{4}; -; \frac{D}{2}, \frac{D+1}{2}; & & \frac{4y^2}{(z + h/2)^2}, \frac{-\alpha^2}{(z + h/2)^2}, \left(\frac{z - h/2}{z + h/2}\right)^2 \\ \frac{E}{2}, \frac{E+1}{2} :: -; -; -; \frac{1}{2}, \frac{5}{4}; \nu + 1; \frac{1}{2} & ; & \end{matrix} \right] \right\} \end{aligned}$$

$$\begin{aligned}
 & + \frac{CD}{E} \left( \frac{z - h/2}{z + h/2} \right) \\
 F^{(3)} & \left[ \begin{array}{c} \frac{C+1}{2}, \frac{C+2}{2} :: \frac{C'}{2}, \frac{C'+1}{2}; -; -; \frac{1}{4}; -; \frac{D+1}{2}, \frac{D+2}{2} ; \\ \frac{E+1}{2}, \frac{E+2}{2} :: -; -; -; \frac{1}{2}, \frac{5}{4}; \nu + 1; \frac{3}{2} \end{array} ; \frac{4y^2}{(z + h/2)^2}, \frac{-\alpha^2}{(z + h/2)^2}, \left( \frac{z - h/2}{z + h/2} \right)^2 \right] \Bigg\} \quad (3.6)
 \end{aligned}$$

(vii) On setting  $\gamma = \frac{3}{2}$ , in equation (2.1) and using a known result [7, p.595(11)] and integral transform [3, p.216(16)], we get

$$\begin{aligned}
 & \frac{\sqrt{\pi}}{2} \Gamma(\delta) y^{1/2-\delta} \int_0^\infty t^{\sigma+1/2-\delta-1} e^{-zt} W_{k,m}(ht) J_\nu(\alpha t) [L_{\delta-3/2}(2yt)] dt \\
 & = \frac{h^{m+\frac{1}{2}} \alpha^\nu \Gamma(C) \Gamma(C')}{2^\nu \Gamma(\nu + 1) \Gamma(E) \left(z + \frac{h}{2}\right)^C} \\
 & \cdot \left\{ F^{(3)} \left[ \begin{array}{c} \frac{C}{2}, \frac{C+1}{2} :: \frac{C'}{2}, \frac{C'+1}{2}; -; -; \beta; -; \frac{D}{2}, \frac{D+1}{2} ; \\ \frac{E}{2}, \frac{E+1}{2} :: -; -; -; \frac{3}{2}, \delta; \nu + 1; \frac{1}{2} \end{array} ; \frac{4y^2}{\left(z + \frac{h}{2}\right)^2}, \frac{-\alpha^2}{\left(z + \frac{h}{2}\right)^2}, \left( \frac{z - \frac{h}{2}}{z + \frac{h}{2}} \right)^2 \right] \right. \\
 & \quad \left. + \frac{CD}{E} \left( \frac{z - h/2}{z + h/2} \right) \right. \\
 & \left. F^{(3)} \left[ \begin{array}{c} \frac{C+1}{2}, \frac{C+2}{2} :: \frac{C'}{2}, \frac{C'+1}{2}; -; -; \beta; -; \frac{D+1}{2}, \frac{D+2}{2} ; \\ \frac{E+1}{2}, \frac{E+2}{2} :: -; -; -; \frac{3}{2}, \delta; \nu + 1; \frac{3}{2} \end{array} ; \frac{4y^2}{(z + h/2)^2}, \frac{-\alpha^2}{(z + h/2)^2}, \left( \frac{z - \frac{h}{2}}{z + \frac{h}{2}} \right)^2 \right] \right\} \quad (3.7)
 \end{aligned}$$

### 4 Conclusion

In the present investigation, we have established on infinite series of three variables involving Whittaker, Bessel and generalized hypergeometric functions into a series of Srivastava’s triple hypergeometric function  $F^{(3)}$ . A number of known and new transformations for Kampe’ de F’erriet function  $F^{(2)}$  and Legendre function are obtained as special cases. The result appears in the paper may be found useful in some areas of mathematical physics and engineering.

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