

# A CUBIC B-SPLINE FINITE ELEMENT METHOD FOR VOLTERRA INTEGRO-DIFFERENTIAL EQUATION

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Communicated by Martin Bohner

MSC 2010 Classifications: Primary 20M99, 13F10; Secondary 13A15, 13M05.

Keywords and phrases: Cubic B-spline method, Volterra Integro-differential equation, priori error estimate, convergence analysis.

*The authors would like to thank the reviewers and editor for their constructive comments and valuable suggestions that improved the quality of our paper.*

*The authors would like to thank of the financial support from Koya University and Imam Ja'afar AL-Sadiq University.*

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**Abstract** This work proposes a numerical solution for the Volterra integro-differential equation (VIDE). The method proposed in this study combines the cubic B-spline method with the finite element method to effectively solve this type of equation. To handle the integral part of the equation. Gauss rules are employed to approximate the integral of the differential equations. These rules ensure the accuracy employed to approximate the integral of differential equations. These rules ensure the accuracy and precision of the results obtained. In addition, the study investigates the coercivity and continuity of the the suggested problems to address challenges in deriving a priori error estimates. Through this analysis, the convergence of the method is proven. To demonstrate the effectiveness of the proposed method, several examples are presented. These examples serve to illustrate the applicability of the method and also allow for a comparison between the approximate results and analytical solutions.

## 1 Introduction

Initial- and boundary-value problems with integro-differential equations (IEDs) are abundant in physical and biological modeling as well as in respective (bio-)engineering applications. In recent years, the finite element method (FEM) has gained significant attention as a numerical technique for approximating solutions to ordinary and partial differential equations that do not involve memory effects [1, 2, 3, 3, 5].

However, the implementation of FEM with cubic B-spline functions for solving VIDE has seen slower progress compared to other methods. As a result, researchers have dedicated considerable effort to developing efficient numerical schemes for approximating solutions to these problems. Numerical solutions for such equations have been explored using various techniques, Galerkin finite element method [6, 7, 8, 9]. Various methods have been applied to solve IEDs. For more details see [10, 11, 12, 13, 14, 15, 16, 17, 18, 19, 20, 21, 22, 23, 24, 25].

The aim of this study is to explore the use of modified cubic B-spline finite element methods to approximate solutions for VIDE. However, there are challenges involved in constructing a linear system that can yield highly accurate results. One of the challenges is integrating the integro parts of the equation. To address this, a quadrature rule can be employed to obtain exact approximations for the integral parts. Another challenge is dealing with Dirichlet boundary conditions using B-spline functions. This requires finding appropriate techniques to incorporate these boundary conditions into the framework of the B-spline finite element method

The present study is prepared as follows: In Section 2, the numerical method of the presented problem was given. Convergence analysis of suggested schemes is proved in Section 3. Some numerical experiments are shown in Section 4. Finally, conclusions are given in Section 5.

## 2 Second-order linear VIDE

Consider the model problem as

$$-z''(s) + b(s)z'(s) + c(s)z(s) = f(s) + \int_0^s k(s, t)z(t)dt, \tag{2.1}$$

with boundary conditions

$$z(a) = z(b) = 0, \quad s, t \in (0, 1), \tag{2.2}$$

where  $b(s), c(s), f(s)$  and  $k(s, t)$  are continuous real-valued functions on the interval  $[0, 1]$ . The weak form of (2.1) can be done multiplying both sides of the (2.1) by a test function  $z \in H_0^1(0, 1)$  and integrating by parts of (2.1), we obtain

$$\begin{aligned} \int_0^1 z'(x)v'(x)ds + \int_0^1 b(s)z'(s)v(s)ds + \int_0^1 c(s)z(s)v(s)ds \\ = \int_0^1 f(s)v(s)ds + \int_0^1 v(s) \int_0^1 k(s, t)z(t)dt ds, \end{aligned}$$

such that

$$a(z, v) = F(v), \quad v \in H_0^1(0, 1), \tag{2.3}$$

where  $a(z, v)$  is a symmetric bilinear form and  $F(v)$  is a linear functional given by:

$$\begin{aligned} a(z, v) &= \int_0^1 z'(s)v'(s)ds + \int_0^1 b(s)z'(s)v(s)ds + \int_0^1 c(s)z(s)v(s)ds \\ &\quad - \int_0^1 v(s) \int_0^s k(s, t)z(t)dt ds \\ F(v) &= \int_0^1 f(s)v(s)ds. \end{aligned}$$

### 2.1 Finite element approximation

To process with finite element approximation, let  $S^h$  be a finite dimensional space such that  $S^h \subset H_0^1(\Omega)$ . To develop the numerical method for approximating the solution of a boundary value problem (2.1), the interval  $[0, 1]$  is partitioned into  $N + 1$  uniformly spaced points  $x_i$  such that  $0 < x_0 < x_1 < \dots < x_{N-1} < x_N = 1$  and  $h = \frac{1}{N+1}$ . Setting

$$S^h = \{v \in S_3 \in C^2[0, 1] : v(0) = v(1) = 0\}, \tag{2.4}$$

where  $S_3$  is the space of all polynomials of degree  $\leq 3$ . We seek to find an approximation  $Z_h \in S^h$  such that

$$a(Z_h, v_h) = F(v_h), \quad v_h \in S^h, \tag{2.5}$$

where

$$\begin{aligned} a(Z_h, v) &= \int_0^1 Z_h'(s)v'(s)ds + \int_0^1 b(s)Z_h'(s)v(s)ds + \int_0^1 c(s)Z_h(s)v(s)ds \\ &\quad - \int_0^1 v(s) \int_0^s k(s, t)Z_h(t)dt ds, \\ F(v_h) &= \int_0^1 f(s)v_h(s)ds. \end{aligned}$$

**2.2 Calculation and assembly of stiffness and mass matrices**

Let  $N$  be a positive integer number and  $a = 0 < s_0 < s_1 < \dots < s_{N-1} < s_N = b$  and  $h = 1/(N + 1)$  be a uniform partition of  $[a, b]$  and  $h = \frac{b-a}{N}$ . The typical third-degree B-spline basis functions defined as

$$B_i^3(s) = \begin{cases} h^{-3}g_1(s - s_{i-2}), & s \in [s_{i-2}, s_{i-1}] \\ g_2\left(\frac{s-s_{i-1}}{h}\right) & s \in [s_{i-1}, s_i] \\ g_2\left(\frac{s_{i+1}-s}{h}\right) & s \in [s_i, s_{i+1}] \\ h^{-3}g_1(s_{i+2} - s), & s \in [s_{i+1}, s_{i+2}] \\ 0, & \text{otherwise} \end{cases} \tag{2.6}$$

where

$$g_1(s) = s^3, \quad g_2(s) = 1 + 3s + 3s^2 - 3s^3, \text{ for } i = -1, 0, \dots, N + 1.$$

For a sufficiently smooth function  $z(s)$  there always exists a unique third-degree spline  $Z_h$  such that

$$Z_h = \sum_{m=-1}^{N+1} \alpha_m B_m^3(s), \tag{2.7}$$

where  $\alpha_i$  are unknown quantities to be determined from (2.6). For the sake of simplicity, using  $z_i = Z_h(s_i, t)$

$$\begin{cases} Z_i = \alpha_{i-1} + 4\alpha_i + \alpha_{i+1} \\ Z'_i = \frac{3}{h} (\alpha_{i-1} - \alpha_{i+1}) \\ Z''_i = \frac{6}{h^2} (\alpha_{i-1} - 2\alpha_i + \alpha_{i+1}). \end{cases} \tag{2.8}$$

Using (2.7) and boundary conditions (2.2), gives

$$\begin{cases} a_1 = -a_{-1} - 4a_0 \\ a_{N+1} = -a_{N-1} - 4a_N. \end{cases} \tag{2.9}$$

Substituting (2.8) in (2.7), we have

$$Z_h = a_0 [B_0(s) - 4B_{-1}(s)] + a_1 [B_1(s) - B_{-1}(s)] + \dots + a_{N-1} [B_{N-1}(s) - B_{N+1}(s)] + a_N [B_N(s) - 4B_N(s)],$$

so that

$$Z_h = a_0 \overline{B}_0(s) + a_1 \overline{B}_1(s) + \dots + a_{N-1} \overline{B}_{N-1}(s) + a_N \overline{B}_N(s),$$

where

$$\begin{cases} \overline{B}_0(s) = B_0(s) - 4B_{-1}(s) \\ \overline{B}_1(s) = B_1(s) - B_{-1}(s) \\ \overline{B}_i(s) = B_i(s), \quad i = 2, 3, \dots, N - 2 \\ \overline{B}_{N-1}(s) = B_{N-1}(s) - B_{N+1}(s) \\ \overline{B}_N(s) = B_N(s) - 4B_{N+1}(s). \end{cases} \tag{2.10}$$

**Table 1.** The values of  $B_i(s)$  and their derivatives.

	$s_i$	$s_{i-1}$	$s_{i+1}$	else
$B_i(s)$	$1/6$	$2/3$	$1/6$	$0$
$B'_i(s)$	$\frac{-1}{2h}$	$0$	$\frac{1}{2h}$	$0$
$B''_i(s)$	$\frac{1}{h^2}$	$\frac{-2}{h^2}$	$\frac{1}{h^2}$	$0$

Inserting (2.10) in (2.5) with choosing  $v_h = \overline{B}_m(s)$ , gives

$$\sum_{i=0}^N \alpha_i \left( \int_0^1 \overline{(B_i)'}(s) (B_m)'(s) + \int_0^1 b(s) \overline{(B_i)'}(s) (B_m)(s) ds + \int_0^1 c(s) \overline{(B_i)'}(s) (B_m)'(s) ds - \int_0^1 \overline{B}_m(s) \int_0^s k(s,t) \overline{B}_i(t) dt ds \right) = \int_0^1 f(s) \overline{B}_m(s) ds,$$

or

$$(A^e + bB^e + cC^e - D^e) \alpha = F^e, \tag{2.11}$$

where

$$A^e = \int_0^1 \overline{(B_i)'}(s) \overline{(B_m)'}(s) ds = \frac{3}{10h} \begin{bmatrix} 80 & 43 & -20 & -1 \\ 43 & 104 & -14 & -24 \\ -20 & -14 & 80 & -15 \\ -1 & -24 & -15 & 80 \end{bmatrix} \tag{2.12}$$

$$B^e = \int_0^1 \overline{(B_i)'}(s) \overline{(B_m)}(s) ds = \frac{1}{20} \begin{bmatrix} 0 & 133 & 52 & 1 \\ -133 & 0 & 244 & 56 \\ -52 & -244 & 0 & 245 \\ -1 & -56 & -245 & 0 \end{bmatrix} \tag{2.13}$$

$$C^e = \int_0^1 \overline{(B_i)}(s) \overline{(B_m)}(s) ds = \frac{h}{40} \begin{bmatrix} 496 & 733 & 116 & 1 \\ 733 & 2296 & 1190 & 120 \\ 116 & 1190 & 2416 & 1191 \\ 1 & 120 & 1191 & 2416 \end{bmatrix} \tag{2.14}$$

are element matrices. The latter integral  $F^e$  and  $D^e$  will be handled numerically using the composite 3-point Gauss quadrature given by:

$$\int_0^1 g(s) ds = \sum_{k=0}^N \int_{s_k}^{s_{k+1}} g(s) ds = \frac{h}{2} \sum_{k=0}^N \int_{-1}^1 g \left( s_k + \frac{h}{2}(1+t) \right) ds \approx \frac{h}{2} \sum_{k=0}^N \sum_{j=0}^3 C_j g(\xi_{k,j}) \tag{2.15}$$

$$\xi_{k,j} = \frac{h}{2} \xi_j + \frac{s_k + s_{k+1}}{2}, \quad k = 0, 1, \dots, N, \quad i = 1, 2, 3.$$

Therefore, using (2.15) we obtain

$$F^e = \int_0^1 f(s) \overline{(B_i)}(s) ds \approx \frac{h}{2} \sum_{k=0}^N \sum_{j=0}^3 C_j f(\xi_{k,j}) \overline{(B_i)}(\xi_{k,j}), \quad i = 1, 2, \dots, N + 1. \tag{2.16}$$

$$D^e = \int_0^1 k(s,t) \overline{(B_i)}(s) dt ds \approx \sum_{k=0}^N \sum_{j=0}^3 C_j K_{k,j} \overline{(B_i)}(\xi_{k,j}), \quad i = 1, 2, \dots, N + 1. \tag{2.17}$$

To find the coefficient  $\alpha_0, \alpha_1, \dots, \alpha_N$ , we assemble all the matrices in (2.11), imply that

$$(A + B + bC + cD) \alpha = F,$$

for each matrix has  $A^e, B^e, C^e$  and  $D^e$  assembled from the element matrices  $A, B, C,$  and  $D$  such that

$$A = \frac{3}{10h}(-1, -44, 14, 344, 14, -44, -1), B = \frac{1}{20}(1, -108, -622, 0, -622, -108, 1), \text{ and}$$

$$C = \frac{3}{10h}(-1, -24, -15, 80, -15, -24, -1).$$

### 3 Convergence analysis

The aim of this section is to present the error analysis theorems for the proposed method. The lemma below will show (2.3) satisfying the Lax–Milgram theorem.

**Lemma 3.1.** (Existence, uniqueness) *Let  $a(z, v)$  be bilinear form defined by (2.3). Suppose that  $N_1 \leq c(s) \leq N_2, Q_1 \leq b(s) \leq Q_2$  and  $0 \leq b'(s) \leq T_2$ , then  $a(z, v)$  is a V-ellipticity, (2.3) has a unique solution.*

$$a(z, z) \geq C_{covr} \|z\|_a^2, \quad \forall z \in H_0^1(0, 1), \tag{3.1}$$

$$a(z, v) \leq C_{cont} \|z\|_a \cdot \|v\|_a, \quad \forall z, v \in H_0^1(0, 1). \tag{3.2}$$

*Proof.* Recalling weak form in (2.3) along with Cauchy–Schwarz inequality, we have

$$a(z, v) \leq \|z(s)\|_a \|v(s)\|_a + Q_2 \|z(s)\|_a \|v(s)\|_{L^2(0,1)} + N_2 \|z(s)\|_{L^2(0,1)} \|v(s)\|_{L^2(0,1)} + KR \|z(s)\|_a \|v(s)\|_a = (1 + Q_2 + N_2 + KR) \|z(s)\|_a \|v(s)\|_a = C_{cont} \|z\|_a \cdot \|v\|_a.$$

Using Poincaré-Friedrichs inequality

$$\|v(s)\|_{L^2(0,1)} \leq C_p \|v(s)\|_{H_0^1(0,1)} = C_p \|v(s)\|_a,$$

and setting

$$C_{cont} = (1 + Q_2 + N_2 + KR), K = \max |K(s, t)|, x \in [0, 1], t \in [0, x], \text{ and } R = \|1\|_{L^2(0,1)}.$$

The second part of (3.2) of above lemma is proved. To prove the V-ellipticity of  $a(z, v)$ , we have

$$\int_0^1 z'(s)z'(s)ds + \int_0^1 c(s)z(s)z(s)ds \geq \int_0^1 (z')^2(s)ds \geq \frac{1}{1+c} \|z\|_{H^1(0,1)}^2, \tag{3.3}$$

and

$$-\int_0^1 z(s) \int_0^s k(s, t)z(t)dsds = -\frac{1}{2} \int_0^1 b'(s)(z(s))^2 ds \geq -\frac{T_2}{2} \int_0^1 (z(s))^2 ds \geq -\frac{T_2}{2} \|z\|_{H^1(0,1)}^2, \tag{3.4}$$

also

$$-\int_0^1 z(s) \int_0^s k(s, t)z(t)dt ds \geq -|\int_0^1 z(s) \int_0^s k(s, t)z(t)dt ds| \tag{3.5}$$

$$\geq -KR \|z\|_{L^2(0,1)}^2 \geq KR \|z\|_{H^1(0,1)}^2. \tag{3.6}$$

$$\tag{3.7}$$

Combining (3.3)-(3.5), we have

$$a(z, z) \geq \left( \frac{1}{1+c} - \frac{T_2}{2} - KR \right) \|z\|_{H^1(0,1)}^2, \tag{3.8}$$

or

$$a(z, z) \geq C_{covr} \|z\|_{H^1(0,1)}^2, \tag{3.9}$$

where  $C_{covr} = \left( \frac{1}{1+c} - \frac{T_2}{2} - KR \right), c$  is a Poincare’s constant. Thus  $a(z, v)$  is a V-elliptic if  $C_{covr} \geq 0$ . Therefore, by the Lax-Milgram theorem and the V-ellipticity of  $a(z, v)$ , equation (2.3) has a unique solution.  $\square$

### 3.1 Error Estimates in the Energy Norm

This section aims to employ Cea’s lemma to derive an optimal error bound for the finite element approximation for a one-dimensional model problem. To start with, the energy norm  $\|v\|_a$  defined by

$$\|z\|_a^2 = \|z'\|^2 + \|z\|^2 + \int_0^s k(s, t)\|z(t)\|^2 dt = \|z\|_{H^1(0,1)}^2.$$

### 3.2 Definition

Let  $H^1(0, 1)$  be a Hilbert space, and suppose  $a$  be a symmetric and  $H_0^1(0, 1)$ -elliptic bilinear form. We define an inner product as follows

$$(\cdot, \cdot) : H^1(0, 1) \times H^1(0, 1) \longrightarrow \mathbb{R},$$

$$(z, v)_a = a(z, v),$$

which is called the inner product energy. Also we define energy norm as follows

$$\|z\|_a^2 = a(z, z).$$

Suppose that  $z$  is the exact solution of the problem and  $Z_h$  be its approximate solution, then, we have

$$a(z, v) = (f(x), v), \quad \forall v \in S^h$$

$$a(Z_h, v_h) = (f(x), v_h), \quad \forall v_h \in S^h.$$

If  $e = z - Z_h$ , then

$$a(e, v_h) = 0, \quad \forall v_h \in S^h. \tag{3.10}$$

The error can be decomposed as

$$e = z - Z_h = z - \Pi_h z + \Pi_h z - Z_h = \epsilon_1 - \epsilon_2,$$

where

$$\epsilon_1 = z - \Pi_h z, \quad \epsilon_2 = \Pi_h z - Z_h,$$

and

$$\Pi_h z : H^1(0, 1) \longrightarrow S^h$$

is the interpolant of the exact solution  $z \in H^1(0, 1)$ . Applying energy norm in above, gives

$$\|e\|_a = \|z - Z_h\|_a \leq \|z - \Pi_h z\|_a + \|\Pi_h z - Z_h\|_a \leq \|\epsilon_1\|_a + \|\epsilon_2\|_a.$$

**Lemma 3.2.** *Let  $S^h \subseteq H_0^1(0, 1)$ , where  $0 \leq h \leq 1$ . For  $z \in H_0^1(0, 1) \cap H^4(0, 1)$ , there exists a constant  $C$  independent of  $h$  such that*

$$\inf_{v_h \in S^h} \|\Pi_h z - Z_h\|_a \leq Ch^3.$$

*Proof.* See [26]. □

**Theorem 3.3.** Assume that  $z$  be the exact solution of (2.3) and  $Z_h$  be the approximate solution of the variational formulation (2.5), then

$$\|z - Z_h\|_a \leq C_1 h^3,$$

where  $C_1 = \frac{CC_{cont}}{C_{covr}} + C$ .

*Proof.*

$$\|e\|_a = z - Z_h = z - \Pi_h z + \Pi_h z - Z_h = \epsilon_1 + \epsilon_2,$$

where  $\epsilon_1 = z - \Pi_h z$ ,  $\epsilon_2 = \Pi_h z - Z_h$ . Plugging  $e = \epsilon_1 + \epsilon_2$  in (3.10), this becomes

$$a(\epsilon_1 + \epsilon_2, v) = a(\epsilon_1, v) + a(\epsilon_2, v) = 0.$$

Testing  $v = \epsilon_2$ , gives

$$a(\epsilon_1, \epsilon_2) + a(\epsilon_2, \epsilon_2) = 0.$$

Therefore

$$a(\epsilon_1, \epsilon_2) = -a(\epsilon_2, \epsilon_2),$$

using the (3.1) and (3.2), imply that

$$C_{\text{covr}} \|\epsilon_2\|_a^2 \leq C_{\text{cont}} \|\epsilon_1\|_a \|\epsilon_2\|_a$$

which leads to

$$\|\epsilon_2\|_a \leq \frac{C_{\text{cont}}}{C_{\text{covr}}} \|\epsilon_1\|_a$$

from Lemma (3.2), gives

$$\|\epsilon_1\|_a = Ch^3.$$

Finally, we obtain

$$\|e\|_a \leq \|\epsilon_1\|_a + \|\epsilon_2\|_a = \frac{CC_{\text{cont}}}{C_{\text{covr}}} h^3 + Ch^3 = C_1 h^3.$$

□

### 4 Numerical Experiments

The section illustrates the performance of a presented method, through an implementation based on Mathematica programming. The error norms are used to measure the error between the numerical and exact solution. The error is defined as

$$E(s) = z(s) - Z_h(s).$$

Then, the pointwise error is

$$\epsilon_A(s) = |E(s_i)|.$$

**Example 4.1.** Given the following first order VIDE

$$z'(s) + z(s) = 1 - s - \frac{2}{3}s^3 - 2 \int_0^s z(t)dt,$$

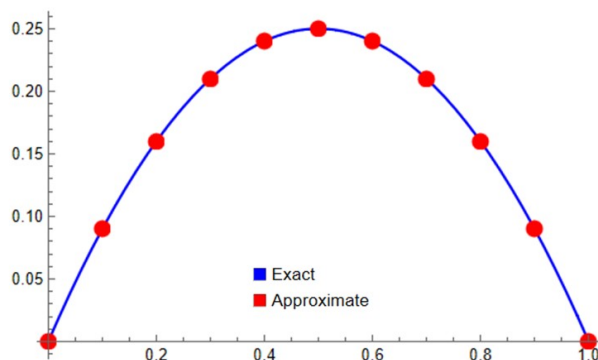
with initial condition

$$z(0) = 0,$$

and exact solution  $z(s) = s - s^2$ .

**Table 2.** Pointwise error  $\epsilon_A(s)$  for Example 4.1.

$s$	$Z_h(s)$	$z(s)$	$\epsilon_A(s)$
0.1	$9.000e^{-02}$	$9.e^{-02}$	$1.544e^{-06}$
0.2	$1.600e^{-02}$	$1.6e^{-01}$	$2.487e^{-06}$
0.3	$2.100e^{-01}$	$2.1e^{-01}$	$7.297e^{-06}$
0.4	$2.399e^{-01}$	$2.4e^{-01}$	$1.603e^{-05}$
0.5	$2.500e^{-01}$	$2.5e^{-01}$	$6.851e^{-06}$
0.6	$2.399e^{-01}$	$2.4e^{-01}$	$1.564e^{-05}$
0.7	$2.100e^{-01}$	$2.1e^{-01}$	$8.723e^{-06}$
0.8	$1.600e^{-01}$	$1.6e^{-01}$	$2.841e^{-06}$
0.9	$9.000e^{-02}$	$9.e^{-02}$	$1.470e^{-06}$



**Figure 1.** Numerical and exact solution of Example 4.1.

**Example 4.2.** Consider the VIDE

$$z''(s) - z(s) = -f(s) - \int_0^s (s-t)z(t)dt, \quad z(0) = z(1) = 0,$$

where

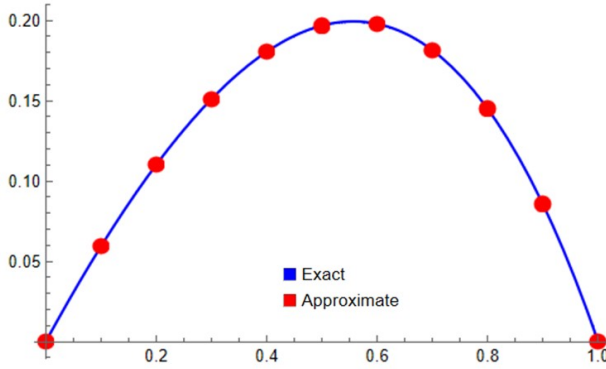
$$f(s) = s - \frac{s^3}{6} - se^{s-1} - e^{-1}(2+s),$$

and the exact solution  $z(s) = s - se^{s-1}$ .

**Table 3.** Pointwise error  $\varepsilon_A(s)$  for Example 4.2.

$s$	$Z_h(s)$	$z(s)$	$\varepsilon_A(s)$
0.1	$5.930e^{-02}$	$5.934e^{-02}$	$3.526e^{-05}$
0.2	$1.100e^{-01}$	$1.101e^{-01}$	$4.026e^{-05}$
0.3	$1.509e^{-01}$	$1.510e^{-01}$	$9.038e^{-05}$
0.4	$1.804e^{-01}$	$1.804e^{-01}$	$6.642e^{-05}$
0.5	$1.966e^{-01}$	$1.967e^{-01}$	$8.194e^{-05}$
0.6	$1.977e^{-01}$	$1.978e^{-01}$	$7.125e^{-05}$
0.7	$1.813e^{-01}$	$1.814e^{-01}$	$9.604e^{-05}$
0.8	$1.449e^{-01}$	$1.450e^{-01}$	$4.870e^{-05}$
0.9	$8.560e^{-02}$	$8.564e^{-02}$	$4.292e^{-05}$





**Figure 2.** Numerical and exact solution of Example 4.2.

**Example 4.3.** Consider the VIDE

$$-z''(s) - 3z(s) = f(s) - \int_0^s \sin(s+t)z(t)dt, \quad z(0) = z(1) = 0,$$

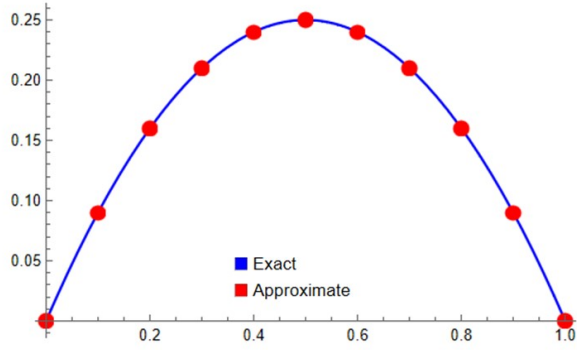
where

$$f(s) = (2 - 3s + 3s^2) + (s^2 - s - 2) \cos(2s) - (2s - 1) \sin(2s) - \sin(2s) + 2 \cos(s),$$

and the exact solution  $z(s) = -s^2 + s$ .

**Table 4.** Pointwise error  $\varepsilon_A(s)$  for Example 4.3.

$s$	$Z_h(s)$	$z(s)$	$\varepsilon_A(s)$
0.1	$8.993e^{-02}$	$9.e^{-02}$	$6.640e^{-05}$
0.2	$1.599e^{-01}$	$1.6e^{-01}$	$8.855e^{-05}$
0.3	$2.098e^{-01}$	$2.1e^{-01}$	$1.633e^{-04}$
0.4	$2.398e^{-01}$	$2.4e^{-01}$	$1.390e^{-04}$
0.5	$2.498e^{-01}$	$2.5e^{-01}$	$1.578e^{-04}$
0.6	$2.398e^{-01}$	$2.4e^{-01}$	$1.378e^{-04}$
0.7	$2.098e^{-01}$	$2.1e^{-01}$	$1.611e^{-04}$
0.8	$1.599e^{-01}$	$1.6e^{-01}$	$8.617e^{-05}$
0.9	$8.993e^{-02}$	$9e^{-02}$	$6.469e^{-05}$



**Figure 3.** Numerical and exact solution of Example 4.3.

**Example 4.4.** Consider the second order VIDE

$$-z''(s) + 4z(s) = -f(s) - \int_0^s stz(t)dt, \quad z(0) = z(1) = 0,$$

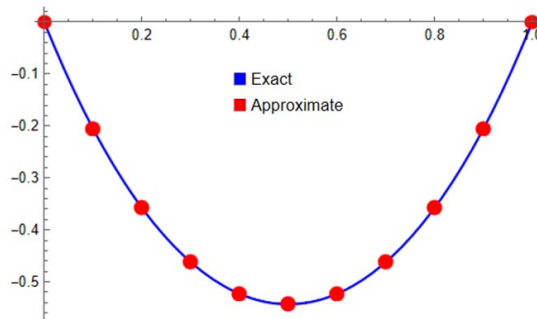
where

$$f(s) = \frac{s^3}{2} \cosh(1) - \frac{s^2}{2} \sinh(2s - 1) + \frac{s}{4} \cosh(2s - 1) - \frac{s}{4} \cosh(1) - 4 \cosh(1),$$

and the exact solution  $z(s) = \cosh(2s - 1) - \cosh(1)$ .

**Table 5.** Pointwise error  $\varepsilon_A(s)$  for Example 4.3.

$s$	$Z_h(s)$	$z(s)$	$\varepsilon_A(s)$
0.1	$8.993e^{-02}$	$9.e^{-02}$	$6.640e^{-05}$
0.2	$1.599e^{-01}$	$1.6e^{-01}$	$8.855e^{-05}$
0.3	$2.098e^{-01}$	$2.1e^{-01}$	$1.633e^{-04}$
0.4	$2.398e^{-01}$	$2.4e^{-01}$	$1.390e^{-04}$
0.5	$2.498e^{-01}$	$2.5e^{-01}$	$1.578e^{-04}$
0.6	$2.398e^{-01}$	$2.4e^{-01}$	$1.378e^{-04}$
0.7	$2.098e^{-01}$	$2.1e^{-01}$	$1.611e^{-04}$
0.8	$1.599e^{-01}$	$1.6e^{-01}$	$8.617e^{-05}$
0.9	$8.993e^{-02}$	$9e^{-02}$	$6.469e^{-05}$



**Figure 4.** Numerical and exact solution of Example 4.4.

This pointwise error norm of  $\varepsilon_A(s)$  is presented in Tables 2 – 5. These tables provide information about the accuracy of the numerical solutions obtained using the cubic B-spline finite element method. Based on the results presented in these tables, it can be concluded that the method is reliable and produces results that are consistent with analytical solutions. This suggests that the numerical method is effective in approximating the solution to the given problem. Furthermore, the physical behavior of both the exact and approximate solutions at different levels of spatial discretization  $h \leq 1$  is depicted graphically in Figures 1 – 4. These figures visually represent how well the numerical solutions capture the true behavior of the system, thereby providing additional evidence of the method's reliability and accuracy.

## 5 Conclusion

This study combines a finite element method with a cubic B-spline function to solve second-order VIDE. The work not only focuses on developing the numerical method but also addresses the existence and uniqueness of solutions and verifies the accuracy of error estimates. To evaluate the precision of the proposed scheme, several test problems are considered, and error norms are calculated for different levels of spatial discretization. The numerical experiments demonstrate that the results obtained from the method are efficient, reliable, fruitful, and powerful in approximating the solutions to the given problems. Moreover, the performance for the method shows good agreement between the exact solution and the approximate solution. However, it would also be interesting to explore the application of a compact finite difference method for solving such problems. For more detailed information, refer to the references [29, 30, 31, 32].

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Received: 2023-12-17

Accepted: 2024-06-06