# ONE PARAMETER FAMILY OF $\mathcal{B}$ - TANGENT DEVELOPABLE SURFACES OF SPACELIKE BIHARMONIC NEW TYPE $\mathcal{B}$ -SLANT HELICES IN $\mathcal{H}^3$

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**Abstract**. In this paper, we study inextensible flows of tangent developable surfaces of biharmonic spacelike new type  $\mathcal{B}$ -slant helices according to Bishop frame in the Lorentzian Heisenberg group  $\mathcal{H}^3$ . We give necessary and sufficient conditions for new type  $\mathcal{B}$ -slant helices to be biharmonic. We characterize one parameter family of  $\mathcal{B}$ -tangent developable surfaces in the Lorentzian Heisenberg group  $\mathcal{H}^3$ . Additionally, we illustrate our results.

## **1** Introduction

A smooth map  $\phi : N \longrightarrow M$  is said to be biharmonic if it is a critical point of the bienergy functional:

$$E_2(\phi) = \int_N \frac{1}{2} \left| \mathcal{T}(\phi) \right|^2 dv_h,$$

where  $\mathcal{T}(\phi) := \operatorname{tr} \nabla^{\phi} d\phi$  is the tension field of  $\phi$ 

The Euler–Lagrange equation of the bienergy is given by  $\mathcal{T}_2(\phi) = 0$ . Here the section  $\mathcal{T}_2(\phi)$  is defined by

$$\mathcal{T}_{2}(\phi) = -\Delta_{\phi}\mathcal{T}(\phi) + \operatorname{tr} R\left(\mathcal{T}(\phi), d\phi\right) d\phi$$

and called the bitension field of  $\phi$ . Non-harmonic biharmonic maps are called proper biharmonic maps.

This study is organised as follows: Firstly, we give necessary and sufficient conditions for new type  $\mathcal{B}$ -slant helices to be biharmonic. We characterize this curves in the Lorentzian Heisenberg group  $\mathcal{H}^3$ . Secondly, we study biharmonic  $\mathcal{B}$ -tangent developable surfaces of spacelike new type  $\mathcal{B}$ -slant helices according to Bishop frame in the Lorentzian Heisenberg group  $\mathcal{H}^3$ . Finally, we illustrate our results.

# **2** The Lorentzian Heisenberg Group $\mathcal{H}^3$

The Heisenberg group  $\text{Heis}^3$  is a Lie group which is diffeomorphic to  $\mathbb{R}^3$  and the group operation is defined as

$$(x,y,z)*(\overline{x},\overline{y},\overline{z})=(x+\overline{x},y+\overline{y},z+\overline{z}-\overline{x}y+x\overline{y}).$$

The identity of the group is (0,0,0) and the inverse of (x, y, z) is given by (-x, -y, -z). The left-invariant Lorentz metric on Heis<sup>3</sup> is

$$g = -dx^2 + dy^2 + (xdy + dz)^2.$$

The following set of left-invariant vector fields forms an orthonormal basis for the corresponding Lie algebra:

$$\left\{\mathbf{e}_1 = \frac{\partial}{\partial z}, \ \mathbf{e}_2 = \frac{\partial}{\partial y} - x\frac{\partial}{\partial z}, \ \mathbf{e}_3 = \frac{\partial}{\partial x}\right\}.$$
 (1)

The characterising properties of this algebra are the following commutation relations, [15]:

$$g(\mathbf{e}_1, \mathbf{e}_1) = g(\mathbf{e}_2, \mathbf{e}_2) = 1, \ g(\mathbf{e}_3, \mathbf{e}_3) = -1.$$

**Proposition 2.1.** For the covariant derivatives of the Levi-Civita connection of the leftinvariant metric g, defined above the following is true:

$$\nabla = \frac{1}{2} \begin{pmatrix} 0 & \mathbf{e}_3 & \mathbf{e}_2 \\ \mathbf{e}_3 & 0 & \mathbf{e}_1 \\ \mathbf{e}_2 & -\mathbf{e}_1 & 0 \end{pmatrix},$$
(2)

where the (i, j)-element in the table above equals  $\nabla_{\mathbf{e}_i} \mathbf{e}_j$  for our basis

 $\{\mathbf{e}_k, k = 1, 2, 3\}.$ 

# 3 Spacelike Biharmonic New Type *B*-Slant Helices with Bishop Frame In The Lorentzian Heisenberg Group *H*<sup>3</sup>

Let  $\gamma : I \longrightarrow \mathcal{H}^3$  be a non geodesic spacelike curve on the Lorentzian Heisenberg group  $\mathcal{H}^3$  parametrized by arc length. Let  $\{\mathbf{t}, \mathbf{n}, \mathbf{b}\}$  be the Frenet frame fields tangent to the Lorentzian Heisenberg group  $\mathcal{H}^3$  along  $\gamma$  defined as follows:

t is the unit vector field  $\gamma'$  tangent to  $\gamma$ , n is the unit vector field in the direction of  $\nabla_t t$  (normal to  $\gamma$ ), and b is chosen so that  $\{t, n, b\}$  is a positively oriented orthonormal basis. Then, we have the following Frenet formulas:

$$\nabla_{\mathbf{t}} \mathbf{t} = \kappa \mathbf{n},$$

$$\nabla_{\mathbf{t}} \mathbf{n} = \kappa \mathbf{t} + \tau \mathbf{b},$$

$$\nabla_{\mathbf{T}} \mathbf{B} = \tau \mathbf{n},$$
(3)

where  $\kappa$  is the curvature of  $\gamma$  and  $\tau$  is its torsion and

$$g(\mathbf{t}, \mathbf{t}) = 1, \ g(\mathbf{n}, \mathbf{n}) = -1, \ g(\mathbf{b}, \mathbf{b}) = 1,$$
$$g(\mathbf{t}, \mathbf{n}) = g(\mathbf{t}, \mathbf{b}) = g(\mathbf{n}, \mathbf{b}) = 0.$$

In the rest of the paper, we suppose everywhere  $\kappa \neq 0$  and  $\tau \neq 0$ .

The Bishop frame or parallel transport frame is an alternative approach to defining a moving frame that is well defined even when the curve has vanishing second derivative. The Bishop frame is expressed as

$$\nabla_{\mathbf{t}} \mathbf{t} = k_1 \mathbf{m}_1 - k_2 \mathbf{m}_2,$$

$$\nabla_{\mathbf{t}} \mathbf{m}_1 = k_1 \mathbf{t},$$

$$\nabla_{\mathbf{t}} \mathbf{m}_2 = k_2 \mathbf{t},$$
(4)

where

$$g(\mathbf{t}, \mathbf{t}) = 1, \ g(\mathbf{m}_1, \mathbf{m}_1) = -1, \ g(\mathbf{m}_2, \mathbf{m}_2) = 1,$$
  
 $g(\mathbf{T}, \mathbf{M}_1) = g(\mathbf{t}, \mathbf{m}_2) = g(\mathbf{m}_1, \mathbf{m}_2) = 0.$ 

Here, we shall call the set  $\{\mathbf{t}, \mathbf{m}_1, \mathbf{m}_2\}$  as Bishop trihedra,  $k_1$  and  $k_2$  as Bishop curvatures. Also,  $\tau(s) = \psi'(s)$  and  $\kappa(s) = \sqrt{|k_2^2 - k_1^2|}$ . Thus, Bishop curvatures are defined by

$$k_1 = \kappa(s) \sinh \psi(s) ,$$
  

$$k_2 = \kappa(s) \cosh \psi(s) .$$

With respect to the orthonormal basis  $\{e_1, e_2, e_3\}$  we can write

$$\mathbf{t} = t^{1}\mathbf{e}_{1} + t^{2}\mathbf{e}_{2} + t^{3}\mathbf{e}_{3},$$
  

$$\mathbf{m}_{1} = m_{1}^{1}\mathbf{e}_{1} + m_{1}^{2}\mathbf{e}_{2} + m_{1}^{3}\mathbf{e}_{3},$$
  

$$\mathbf{m}_{2} = m_{2}^{1}\mathbf{e}_{1} + m_{2}^{2}\mathbf{e}_{2} + m_{2}^{3}\mathbf{e}_{3}.$$
(5)

**Theorem 3.1.**  $\gamma : I \longrightarrow \mathcal{H}^3$  is a spacelike biharmonic curve with Bishop frame if and only if

$$k_{1}^{2} - k_{2}^{2} = \text{constant} = C \neq 0,$$
  

$$k_{1}^{\prime\prime} + \left[k_{1}^{2} - k_{2}^{2}\right] k_{1} = -k_{1} \left[1 + \left(m_{2}^{1}\right)^{2}\right] + k_{2}m_{1}^{1}m_{2}^{1},$$

$$k_{2}^{\prime\prime} + \left[k_{1}^{2} - k_{2}^{2}\right] k_{2} = -k_{1}m_{1}^{1}m_{2}^{1} - k_{2} \left[-1 + \left(m_{1}^{1}\right)^{2}\right].$$
(6)

# 4 B-Tangent Developable Surfaces of Spacelike Biharmonic New Type B-Slant Helices with Bishop Frame In The Lorentzian Heisenberg Group H<sup>3</sup>

To separate a tangent developable according to Bishop frame from that of Frenet- Serret frame, in the rest of the paper, we shall use notation for this surface as  $\mathcal{B}$ -tangent developable.

The purpose of this section is to study  $\mathcal{B}$ -tangent developable of biharmonic spacelike new type  $\mathcal{B}$ -slant helix in  $\mathcal{H}^3$ . The  $\mathcal{B}$ -tangent developable of  $\gamma$  is a ruled surface

$$\mathcal{O}_{new}\left(s,u\right) = \gamma\left(s\right) + u\gamma'\left(s\right). \tag{7}$$

**Definition 4.1.** A surface evolution  $\mathcal{O}_{new}(s, u, t)$  and its flow  $\frac{\partial \mathcal{O}_{new}}{\partial t}$  are said to be inextensible if its first fundamental form  $\{\mathbf{E}, \mathbf{F}, \mathbf{G}\}$  satisfies

$$\frac{\partial \mathbf{E}}{\partial t} = \frac{\partial \mathbf{F}}{\partial t} = \frac{\partial \mathbf{G}}{\partial t} = 0.$$
(8)

**Definition 4.2.** We can define the following one-parameter family of developable ruled surface

$$\mathcal{O}_{new}\left(s, u, t\right) = \gamma\left(s, t\right) + u\gamma'\left(s, t\right).$$
(9)

Hence, we have the following theorem.

**Theorem 4.3.** Let  $\mathcal{O}_{new}$  be one-parameter family of the  $\mathcal{B}$ -tangent developable of a unit speed non-geodesic biharmonic new type  $\mathcal{B}$ -slant helix. Then  $\frac{\partial \mathcal{O}_{new}}{\partial t}$  is inextensible if and only if

$$\frac{\partial}{\partial t} (\sin \mathfrak{Q}(t) - uk_2(t) \cos \mathfrak{Q}(t))^2 + \frac{\partial}{\partial t} (\cos \mathfrak{Q}(t) \cosh [\mathcal{C}_0(t) s + \mathcal{C}_1(t)] 
+ uk_1(t) \sinh [\mathcal{C}_0(t) s + \mathcal{C}_1(t)] - uk_2(t) \sin \mathfrak{Q}(t) \cosh [\mathcal{C}_0(t) s + \mathcal{C}_1(t)])^2 
+ \frac{\partial}{\partial t} (\cos \mathfrak{Q}(t) \sinh [\mathcal{C}_0(t) s + \mathcal{C}_1(t)] + uk_1(t) \cosh [\mathcal{C}_0(t) s + \mathcal{C}_1(t)]$$
(10)  

$$- uk_2(t) \sin \mathfrak{Q}(t) \sinh [\mathcal{C}_0(t) s + \mathcal{C}_1(t)])^2 = 0,$$

where  $C_0, C_1$  are smooth functions of time.

**Proof.** Assume that  $\mathcal{O}_{new}(s, u, t)$  be a one-parameter family of the  $\mathcal{B}$ -tangent developable of a unit speed non-geodesic biharmonic new type  $\mathcal{B}$ -slant helix.

From our assumption, we get the following equation

$$\mathbf{m}_{2} = \cos \mathfrak{Q}(t) \mathbf{e}_{1} + \sin \mathfrak{Q}(t) \cosh \left[\mathcal{C}_{0}(t) s + \mathcal{C}_{1}(t)\right] \mathbf{e}_{2} + \sin \mathfrak{Q}(t) \sinh \left[\mathcal{C}_{0}(t) s + \mathcal{C}_{1}(t)\right] \mathbf{e}_{3}.$$
(11)

where  $C_0, C_1$  are smooth functions of time.

On the other hand, using Bishop formulas Eq.(4) and Eq.(1), we have

$$\mathbf{m}_{1} = \sinh\left[\mathcal{C}_{0}\left(t\right)s + \mathcal{C}_{1}\left(t\right)\right]\mathbf{e}_{2} + \cosh\left[\mathcal{C}_{0}\left(t\right)s + \mathcal{C}_{1}\left(t\right)\right]\mathbf{e}_{3}.$$
(12)

Using above equation and Eq.(11), we get

$$\mathbf{t} = \sin \mathfrak{Q}(t) \mathbf{e}_1 + \cos \mathfrak{Q}(t) \cosh \left[\mathcal{C}_0(t) s + \mathcal{C}_1(t)\right] \mathbf{e}_2 + \cos \mathfrak{Q}(t) \sinh \left[\mathcal{C}_0(t) s + \mathcal{C}_1(t)\right] \mathbf{e}_3.$$
(13)

Furthermore, we have the natural frame  $\{(\mathcal{O}_{new})_s, (\mathcal{O}_{new})_u\}$  given by

$$\left(\mathcal{O}_{new}\right)_{s} = \left(\sin\mathfrak{Q}\left(t\right) - uk_{2}\left(t\right)\cos\mathfrak{Q}\left(t\right)\right)\mathbf{e}_{1} + \left(\cos\mathfrak{Q}\left(t\right)\cosh\left[\mathcal{C}_{0}\left(t\right)s + \mathcal{C}_{1}\left(t\right)\right]\right) + (14)$$

$$\begin{aligned} uk_1(t) \sinh\left[\mathcal{C}_0(t) \, s \, + \mathcal{C}_1(t)\right] &- uk_2(t) \sin \mathfrak{Q}(t) \cosh\left[\mathcal{C}_0(t) \, s + \mathcal{C}_1(t)\right]\right) \mathbf{e}_2 \, + \\ &\left(\cos \mathfrak{Q}(t) \sinh\left[\mathcal{C}_0(t) \, s + \mathcal{C}_1(t)\right] + uk_1(t) \cosh\left[\mathcal{C}_0(t) \, s + \mathcal{C}_1(t)\right] \, - \\ & uk_2(t) \sin \mathfrak{Q}(t) \sinh\left[\mathcal{C}_0(t) \, s + \mathcal{C}_1(t)\right]\right) \mathbf{e}_3, \end{aligned}$$

and

 $\left(\mathcal{O}_{new}\right)_{u} = \sin\mathfrak{Q}\left(t\right)\mathbf{e}_{1} + \cos\mathfrak{Q}\left(t\right)\cosh\left[\mathcal{C}_{0}\left(t\right)s + \mathcal{C}_{1}\left(t\right)\right]\mathbf{e}_{2} + \cos\mathfrak{Q}\left(t\right)\sinh\left[\mathcal{C}_{0}\left(t\right)s + \mathcal{C}_{1}\left(t\right)\right]\mathbf{e}_{3}.$ 

The components of the first fundamental form are

$$\frac{\partial \mathbf{E}}{\partial t} = \frac{\partial}{\partial t} g((\mathcal{O}_{new})_s, (\mathcal{O}_{new})_s) = \frac{\partial}{\partial t} (\sin \mathfrak{Q}(t) - uk_2(t) \cos \mathfrak{Q}(t))^2 
+ \frac{\partial}{\partial t} (\cos \mathfrak{Q}(t) \cosh [\mathcal{C}_0(t) s + \mathcal{C}_1(t)] + uk_1(t) \sinh [\mathcal{C}_0(t) s + \mathcal{C}_1(t)] 
- uk_2(t) \sin \mathfrak{Q}(t) \cosh [\mathcal{C}_0(t) s + \mathcal{C}_1(t)])^2 
+ \frac{\partial}{\partial t} (\cos \mathfrak{Q}(t) \sinh [\mathcal{C}_0(t) s + \mathcal{C}_1(t)] + uk_1(t) \cosh [\mathcal{C}_0(t) s + \mathcal{C}_1(t)] 
- uk_2(t) \sin \mathfrak{Q}(t) \sinh [\mathcal{C}_0(t) s + \mathcal{C}_1(t)])^2, 
\frac{\partial \mathbf{F}}{\partial t} = 0,$$
(15)

$$\frac{\partial \mathbf{G}}{\partial t} = 0.$$

Hence,  $\frac{\partial O_{new}}{\partial t}$  is inextensible if and only if Eq.(10) is satisfied. This concludes the proof of theorem.

**Theorem 4.4.** Let  $\mathcal{O}_{new}$  be one-parameter family of the  $\mathcal{B}$ -tangent developable surface of a unit speed non-geodesic biharmonic new type  $\mathcal{B}$ -slant helix. Then, the parametric equations of this family are given by

$$\begin{split} \boldsymbol{x}_{\mathcal{O}_{new}}\left(s, u, t\right) &= \frac{1}{\mathcal{C}_{0}\left(t\right)} \cos \mathfrak{Q}\left(t\right) \cosh\left[\mathcal{C}_{0}\left(t\right)s + \mathcal{C}_{1}\left(t\right)\right] \\ &+ u \cos \mathfrak{Q}\left(t\right) \sinh\left[\mathcal{C}_{0}\left(t\right)s + \mathcal{C}_{1}\left(t\right)\right] + \mathcal{C}_{2}\left(t\right), \\ \boldsymbol{y}_{\mathcal{O}_{new}}\left(s, u, t\right) &= \frac{1}{\mathcal{C}_{0}\left(t\right)} \cos \mathfrak{Q}\left(t\right) \sinh\left[\mathcal{C}_{0}\left(t\right)s + \mathcal{C}_{1}\left(t\right)\right] \\ &+ u \cos \mathfrak{Q}\left(t\right) \cosh\left[\mathcal{C}_{0}\left(t\right)s + \mathcal{C}_{1}\left(t\right)\right] + \mathcal{C}_{3}\left(t\right), \end{split}$$

$$\boldsymbol{z}_{\mathcal{O}_{new}}\left(s, u, t\right) = \sin \mathfrak{Q}\left(t\right) s - \frac{\mathcal{C}_{2}\left(t\right)}{\mathcal{C}_{0}\left(t\right)} \cos \mathfrak{Q}\left(t\right) \sinh \left[\mathcal{C}_{0}\left(t\right) s + \mathcal{C}_{1}\left(t\right)\right]$$
(16)

$$-\frac{1}{4\mathcal{C}_{0}}\cos^{2}\mathfrak{Q}\left(\mathfrak{t}\right)\left(2[\mathcal{C}_{0}\left(t\right)s+\mathcal{C}_{1}\left(t\right)\right]$$

$$+ \sinh 2[\mathcal{C}_{0}(t) s + \mathcal{C}_{1}(t)]) + u \sin \mathfrak{Q}(t) - \mathfrak{u}(\frac{1}{\mathcal{C}_{0}(t)} \cos \mathfrak{Q}(t) \cosh [\mathcal{C}_{0}(t) s + \mathcal{C}_{1}(t)] + \mathcal{C}_{2}(t)) \cos \mathfrak{Q}(t) \cosh [\mathcal{C}_{0}(t) s + \mathcal{C}_{1}(t)] + \mathcal{C}_{4}(t),$$

where  $C_1, C_2, C_3, C_4, C_0$  are smooth functions of time.

**Proof.** We assume that  $\gamma$  is a unit speed new type  $\mathcal{B}$ -slant helix. Substituting Eq.(1) to Eq.(13), we have

$$\mathbf{t} = (\cos \mathfrak{Q}(t) \sinh [\mathcal{C}_0(t) s + \mathcal{C}_1(t)], \cos \mathfrak{Q}(t) \cosh [\mathcal{C}_0(t) s + \mathcal{C}_1(t)], \sin \mathfrak{Q}(t)$$
(17)

 $-\left(\frac{1}{\mathcal{C}_{0}\left(t\right)}\cos\mathfrak{Q}\left(t\right)\cosh\left[\mathcal{C}_{0}\left(t\right)s+\mathcal{C}_{1}\left(t\right)\right]+\mathcal{C}_{2}\left(t\right)\right)\cos\mathfrak{Q}\left(t\right)\cosh\left[\mathcal{C}_{0}\left(t\right)s+\mathcal{C}_{1}\left(t\right)\right]\right),$ Substituting this into the Eq.(15), we have Eq.(16). Thus, the proof is completed.

We can use Mathematica in above theorem, yields



Figure 1.



Figure 2.

**Fig. 1,2:** The equation (16) is illustrated colour Red, Blue, Purple, Orange, Magenta, Cyan, Yellow, Green at the time t = 1, t = 1.2, t = 1.4, t = 1.6, t = 1.8, t = 2, t = 2.2, t = 2.4, respectively.

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